

ANALYSIS OF THE AVALOS–TRIGGIANI PROBLEM FOR THE LINEAR OSKOLKOV SYSTEM OF THE HIGHEST ORDER AND A SYSTEM OF WAVE EQUATIONS

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The Avalos–Triggiani problem for a system of wave equations and a linear Oskolkov system of the highest order is investigated. The mathematical model contains a linear Oskolkov system describing the flow of an incompressible viscoelastic Kelvin–Voigt fluid of the highest order, and a wave vector equation corresponding to some structure immersed in the specified fluid. Based on the method proposed by the authors of this problem, the theorem of the existence of the unique solution to the Avalos–Triggiani problem for the indicated systems is proved.

Keywords: Avalos–Triggiani problem; incompressible viscoelastic fluid; linear Oskolkov systems.

Introduction

Let Ω be a bounded domain in \mathbb{R}^n , $n = 2, 3$, with sufficiently smooth boundary $\partial\Omega$. Let $u = \text{col}(u_1, u_2, \dots, u_n)$ be a n –dimensional velocity vector $n = 2, 3$, the scalar function p be a pressure, and the vector $w = \text{col}(w_1, w_2, \dots, w_n)$ be a vector of displacement of a body, which occupies the domain Ω_s , and is immersed in a fluid occupying the domain Ω_f . Therefore, $\Omega = \Omega_s \cup \Omega_f$, $\overline{\Omega}_s \cap \overline{\Omega}_f = \partial\Omega_s \equiv \Gamma_s$ is the common boundary of Ω_s , and Ω_f . Let us denote the outer boundary of Ω_f by Γ_f (see Fig. 1). Our goal is to investigate the Avalos–Triggiani problem [1, 2] for the case when the fluid in Ω_f is an incompressible viscoelastic Kelvin–Voigt fluid of the highest order K ($K = n_1 + \dots + n_M$) [3]. The considered mathematical model is determined by the system

$$(1 - \varkappa \nabla^2)u_t - \nu \nabla^2 u + (u \cdot \nabla)u - \sum_{m=1}^M \sum_{s=0}^{n_m-1} A_{m,s} \nabla^2 \mathbf{w}_{m,s} + \nabla p = 0 \quad (1)$$

$$\forall(t, x) \in (0, T] \times \Omega_f \equiv \Omega_{Tf},$$

$$\frac{\partial \mathbf{w}_{m,0}}{\partial t} = u + \alpha_m \mathbf{w}_{m,0}, \quad \alpha_m \in \mathbb{R}_-, \quad m = \overline{1, M} \quad \forall(t, x) \in \Omega_{Tf}, \quad (2)$$

$$\frac{\partial \mathbf{w}_{m,s}}{\partial t} = s \mathbf{w}_{m,s-1} + \alpha_m \mathbf{w}_{m,s}, \quad s = \overline{1, n_m - 1} \quad \forall(t, x) \in \Omega_{Tf}, \quad (3)$$

$$\nabla \cdot u = 0, \quad \forall(t, x) \in \Omega_{Tf}, \quad (4)$$

$$w_{tt} - \nabla^2 w + w = 0 \quad \forall(t, x) \in (0, T] \times \Omega_s \equiv \Omega_{Ts} \quad (5)$$

with the boundary value conditions

$$u|_{\Gamma_f} \equiv 0, \quad \forall(t, x) \in (0, T] \times \Gamma_f \equiv \Gamma_{Tf}, \quad (6)$$

$$\mathbf{w}_{m,s}|_{\Gamma_f} \equiv 0, \quad \forall(t, x) \in \Gamma_{Tf}, \quad (7)$$

$$u \equiv w_t, \quad \forall(t, x) \in (0, T] \times \Gamma_s \equiv \Gamma_{Ts}, \quad (8)$$

$$\frac{\partial u}{\partial \nu} - \frac{\partial w}{\partial \nu} = p\nu \quad \forall(t, x) \in \Gamma_{Ts} \quad (9)$$

and the initial value condition

$$\begin{aligned} (w(0, \cdot), w_t(0, \cdot), \mathbf{w}_{1,0}(0, \cdot), \dots, \mathbf{w}_{M,n_m-1}(0, \cdot), u(0, \cdot)) = \\ = (w_0, w_1, \mathbf{w}_{1,0}^0, \dots, \mathbf{w}_{M,n_m-1}^0) \in \mathbf{H}, \end{aligned} \quad (10)$$

where $\mathbf{H} = (H^1(\Omega_s))^n \times (L^2(\Omega_s))^n \times \mathcal{H}_{1,0} \times \dots \times \mathcal{H}_{M,n_m-1} \times \mathcal{H}_f$ and $\mathcal{H}_{m,s} = (L^2(\Omega_s))^n$, $m = \overline{1, M}, s = \overline{1, n_m - 1}$, $\mathcal{H}_f = \{f \in (L^2(\Omega_f))^n : \nabla \cdot f = 0 \text{ in } \Omega_f \text{ and } [f \cdot \nu]|_{\Gamma_f} = 0\}$.

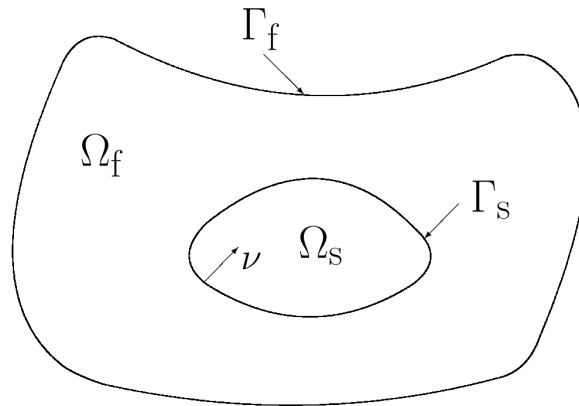


Fig. Physical model

In system (1), the parameters κ and μ characterize the elastic and viscous properties of the fluid, respectively, the parameters $A_{m,s}$ determine the time of pressure retardation (delay), ν is a unit normal vector. In the case of a zero-order Oskolkov system, i.e. $K = 0$, and $\kappa = 0$, problem (1) – (9) was investigated in [1, 2], and for $K = 0$ and $\kappa \neq 0$ – in [4, 5]. The results of this work generalize the results of [6, 7] to the case of the Kelvin–Voigt model of the highest order.

1. Reduction to the Cauchy problem

Following [1, 2], we assume that $p(t)$ satisfies the following elliptic problem:

$$\begin{aligned} \Delta p &= 0 \quad \text{in } \Omega_{Tf}, \\ p &= \frac{\partial u}{\partial \nu} \cdot \nu - \frac{\partial w}{\partial \nu} \cdot \nu \quad \text{on } \Gamma_{Ts}, \\ \frac{\partial p}{\partial \nu} &= \Delta u \cdot \nu \quad \text{on } \Gamma_{Tf}. \end{aligned} \quad (11)$$

Then the pressure p can be represented as follows:

$$p(t) = D_s \left\{ \left(\frac{\partial u(t)}{\partial \nu} \cdot \nu - \frac{\partial w(t)}{\partial \nu} \cdot \nu \right)_{\Gamma_{Ts}} \right\} + N_f((\Delta u(t) \cdot \nu)_{\Gamma_{Tf}}) \quad \text{in } \Omega_{Tf};$$

where the Dirichlet map D_s is defined by the relations

$$h = D_s(g) \Leftrightarrow \begin{cases} \Delta h = 0 & \text{in } \Omega_f, \\ h = g & \text{on } \Gamma_s, \\ \frac{\partial h}{\partial \nu} = 0 & \text{on } \Gamma_f, \end{cases}$$

and the Neumann map N_f is defined by the relations

$$h = N_f(g) \Leftrightarrow \begin{cases} \Delta h = 0 & \text{in } \Omega_f, \\ h = 0 & \text{on } \Gamma_s, \\ \frac{\partial h}{\partial \nu} = g & \text{on } \Gamma_f. \end{cases}$$

Then original system (1) – (5), which describes the interaction of the fluid and the body immersed in the fluid, takes the form

$$(1 - \varkappa \nabla^2)u_t - \nu \nabla^2 u + (u \cdot \nabla)u - \sum_{m=1}^M \sum_{s=0}^{n_m-1} A_{m,s} \nabla^2 \mathbf{w}_{m,s} + G_1 w + G_2 u = 0 \quad (12)$$

$$\forall(t, x) \in (0, T] \times \Omega_f \equiv \Omega_{Tf},$$

$$\frac{\partial \mathbf{w}_{m,0}}{\partial t} = u + \alpha_m \mathbf{w}_{m,0}, \quad \alpha_m \in \mathbb{R}_-, \quad A_{m,s} \in \mathbb{R}_+, \quad m = \overline{1, M} \quad \forall(t, x) \in \Omega_{Tf}, \quad (13)$$

$$\frac{\partial \mathbf{w}_{m,s}}{\partial t} = s \mathbf{w}_{m,s-1} + \alpha_m \mathbf{w}_{m,s}, \quad s = \overline{1, n_m - 1} \quad \forall(t, x) \in \Omega_{Tf}, \quad (14)$$

$$\nabla \cdot u = 0, \quad (15)$$

$$w_{tt} - \nabla^2 w + w = 0 \quad \forall(t, x) \in \Omega_{Ts} \quad (16)$$

with the boundary value conditions

$$u|_{\Gamma_f} \equiv 0, \quad \forall(t, x) \in \Gamma_{Tf}, \quad (17)$$

$$\mathbf{w}_{m,s}|_{\Gamma_f} \equiv 0, \quad \forall(t, x) \in \Gamma_{Tf}, \quad (18)$$

$$u \equiv w_t, \quad \forall(t, x) \in \Gamma_{Ts}, \quad (19)$$

where

$$G_1 w \equiv \nabla \left\{ D_s \left\{ \left(\frac{\partial w(t)}{\partial \nu} \cdot \nu \right)_{\Gamma_{Ts}} \right\} \right\} \quad \text{in } \Omega_{Tf},$$

$$G_2 u \equiv -\nabla \left\{ D_s \left\{ \left(\frac{\partial u(t)}{\partial \nu} \cdot \nu \right)_{\Gamma_{Ts}} \right\} + N_f((\Delta u(t) \cdot \nu)_{\Gamma_{Tf}}) \right\} \quad \text{in } \Omega_{Tf}.$$

Based on the corresponding results for the operators L and M [8], problem (12) – (19), in which pressure is excluded, will be written as an abstract Cauchy problem

$$L\dot{v} = Mv, \quad v(0) = v_0, \quad (20)$$

Here $v = \overline{\text{col}(w, w_t, \mathbf{w}_{1,0}, \dots, \mathbf{w}_{M,0}, \mathbf{w}_{1,1}, \dots, \mathbf{w}_{1,l_1}, \dots, \mathbf{w}_{M,1}, \dots, \mathbf{w}_{M,l_M}, u)}$, where $l_m = n_m - 1, m = \overline{1, M}$.

We study the problem (20) using the results obtained in [9–12].

Lemma 1. *Let $\kappa \in \mathbb{R}, \mu \in \mathbb{R}_+$, the operators L and M be linear continuous operators from \mathbf{G} to \mathbf{H} ($L, M \in \mathcal{L}(\mathbf{G}, \mathbf{H})$), then there exists $L^{-1} \in \mathcal{L}(\mathbf{H})$. Here is the space $\mathbf{G} = \overline{(H^2(\Omega_s))^n \times (H^2(\Omega_s))^n \times \mathcal{G}_{1,0} \times \dots \times \mathcal{G}_{M,n_m-1} \times \mathcal{G}_f}$, where $\mathcal{G}_{m,s} = (H^2(\Omega_s))^n, m = \overline{1, M}, s = \overline{1, n_m - 1}$, \mathcal{G}_f is closure according to the norm of the space $(H^2(\Omega_s))^n$ of the space of infinitely differentiable solenoid functions such that (17) – (19) are fulfilled.*

Theorem 1. *For any $\kappa \in \mathbb{R}, \mu \in \mathbb{R}_+$ and $v_0 \in \mathbf{G}$, there is the unique solution to the problem (20) $v \in C^\infty((0, T], \mathbf{G})$*

In conclusion, we note that the corresponding stochastic models can also be considered using the approach outlined in [13–15].

Acknowledgments. *The work was carried out within the framework of solving problems for the development of the laboratory of Differential Equations and Mathematical Physics of Yaroslav-the-Wise Novgorod State University. The authors express their gratitude to Professor G.A. Sviridyuk for his attention to the work and discussion of the results.*

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Received March 5, 2024

УДК 517.9

DOI: 10.14529/mmp240209

АНАЛИЗ ЗАДАЧИ АВАЛОС – ТРИДЖИАНИ ДЛЯ ЛИНЕЙНОЙ СИСТЕМЫ ОСКОЛКОВА ВЫСШЕГО ПОРЯДКА И СИСТЕМЫ ВОЛНОВЫХ УРАВНЕНИЙ

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В работе исследована задача Авалос – Триджиани для системы волновых уравнений и линейной системы Осколкова высшего порядка. Математическая модель содержит линейную систему Осколкова, описывающую течение несжимаемой вязкоупругой жидкости Кельвина – Фойгта высшего порядка, и волновое векторное уравнение, соответствующее некоторой структуре, погруженной в указанную жидкость. На основе метода, предложенного авторами задачи, доказана теорема существования единственного решения задачи Авалос – Триджиани для указанных систем.

Ключевые слова: задача Авалос – Триджиани, несжимаемая вязкоупругая жидкость, линейные системы Осколкова.

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Поступила в редакцию 5 марта 2024 г.